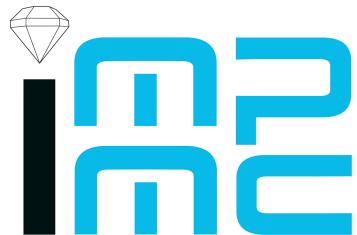


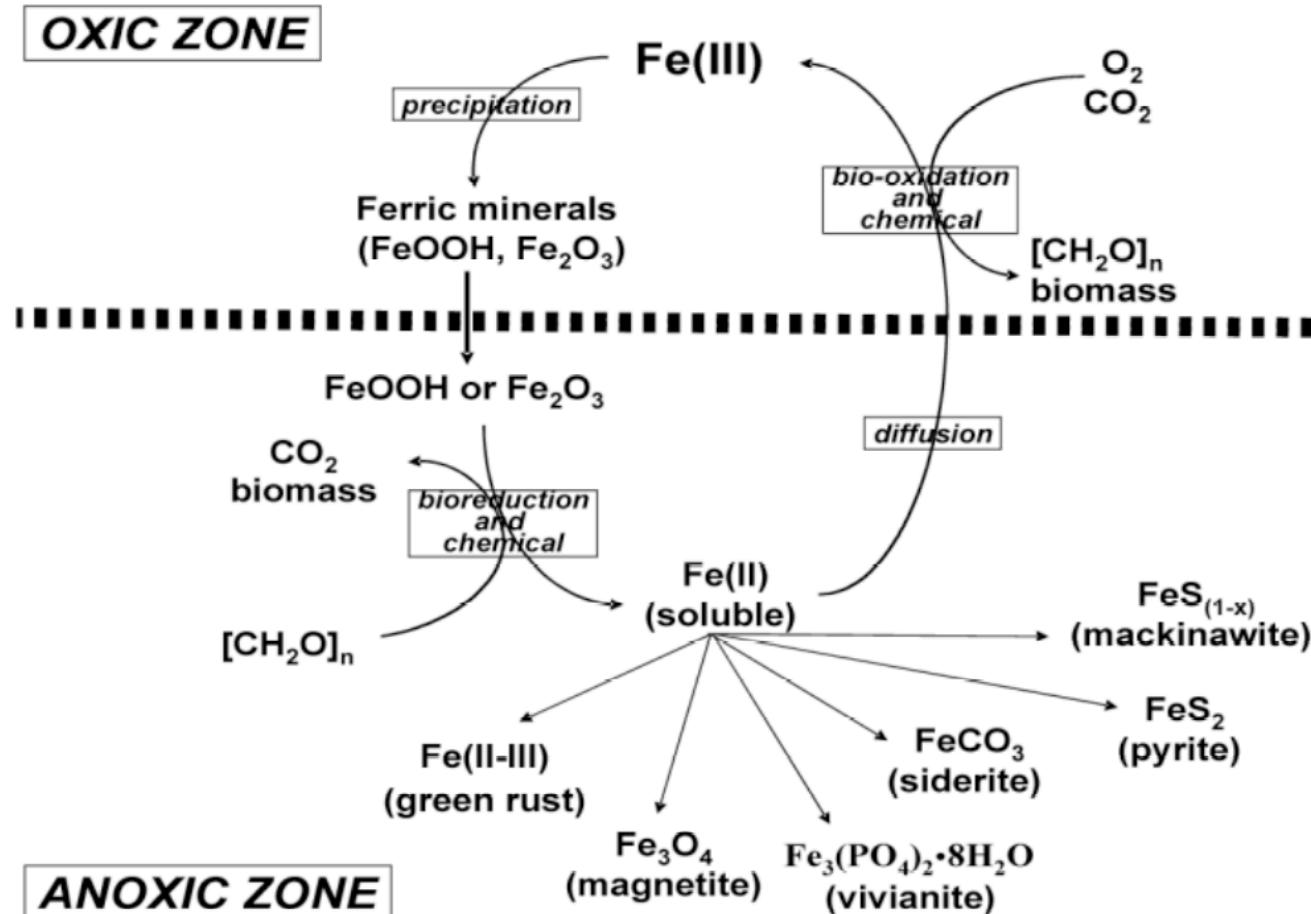
*Theoretical investigation of vibrational and  
isotopic properties of iron (oxyhydr)oxides*

Marc Blanchard & Etienne Balan



## Iron biogeochemistry

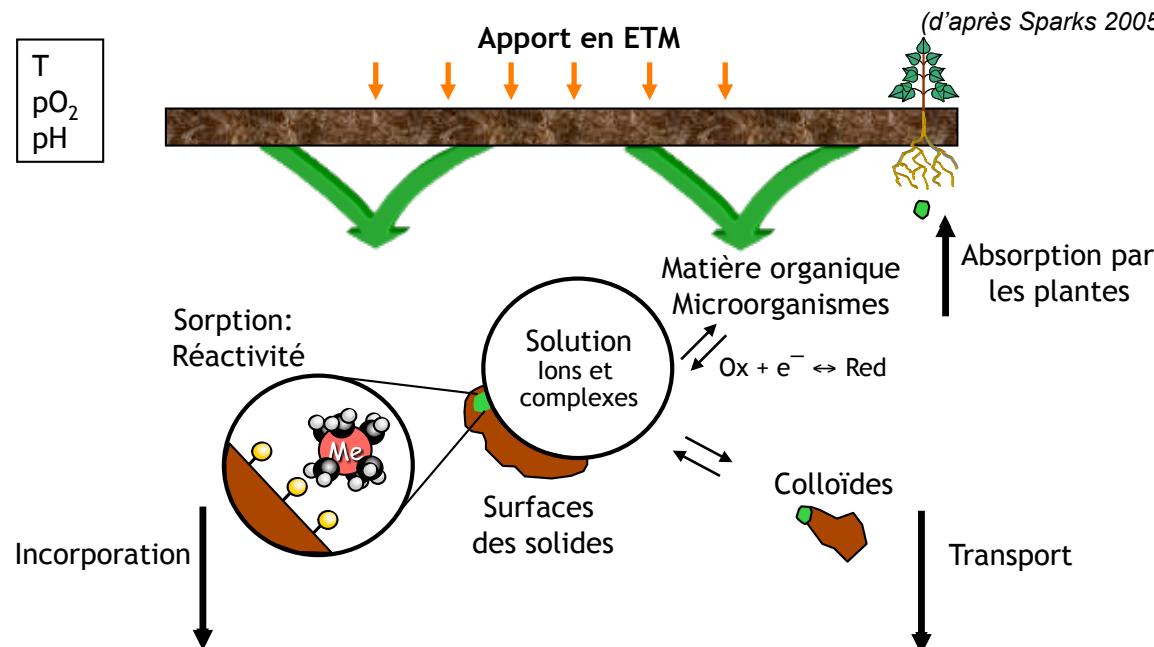
- Iron is the 4<sup>th</sup> most abundant element in the continental crust (5.1 wt%)
- 2 oxidation states (II, III)



Cornell RM, Schwertman U (2003) *The iron oxides*. Wiley-VCH, pp. 664

Chemical reactions at the interfaces between iron (oxyhydr)oxides and water play an important role in numerous natural processes:

- Biological availability (essential nutrient)
- Redox processes (weathering, soil formation, BIF)
- Fate and transport of environmental contaminants

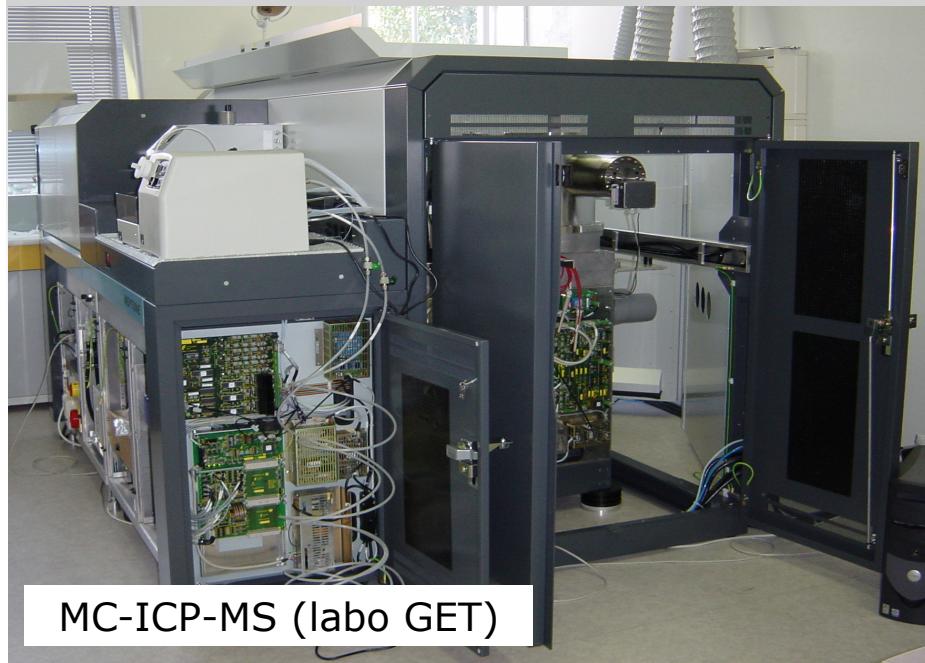


Laterite, Brazil (IMPMC)



Acid mine drainage, Ohio  
(Cornell and Schwertmann 2003)

## Iron isotope geochemistry

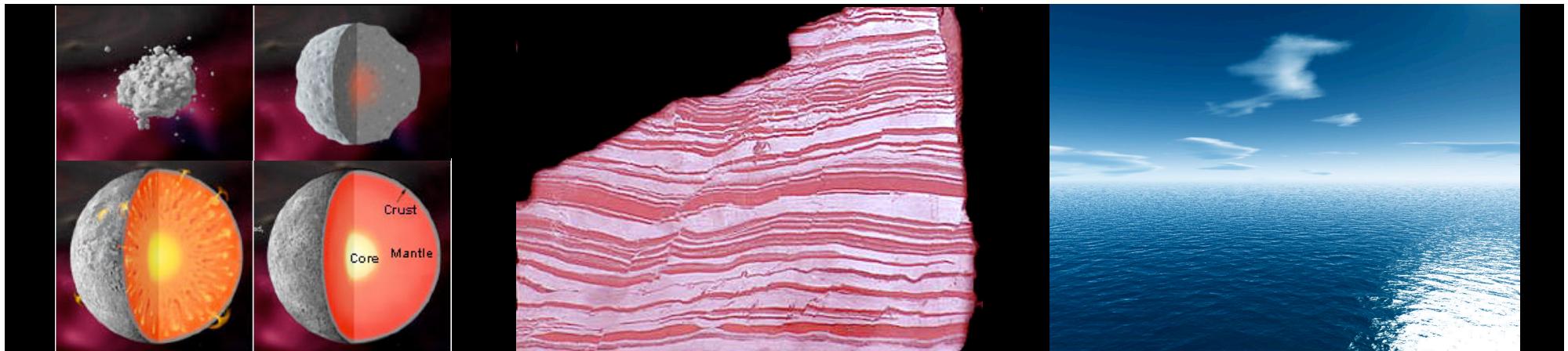


MC-ICP-MS (labo GET)

$^{54}\text{Fe}$  (5.84 %),  $^{56}\text{Fe}$  (91.76 %),  
 $^{57}\text{Fe}$  (2.12 %),  $^{58}\text{Fe}$  (0.28 %)

⇒ Proxy complementary to traditional  
stable isotopes

The most important Fe isotope fractionations occur in low-T environments and with redox gradients.



Planetary formation

Banded Iron Formation

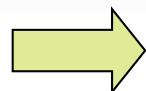
Redox evolution of the ocean

An essential basis for interpreting isotopic compositions in natural samples

is to know the equilibrium isotopic fractionation factors

Mass-dependent equilibrium isotopic fractionation is controlled by

vibrational properties



Need a good description of the vibrational properties (IR, Raman)

For iron (oxyhydr)oxide, occurring most of the time as finely divided minerals,

IR spectroscopy is a technique of choice for the identification, for tracking

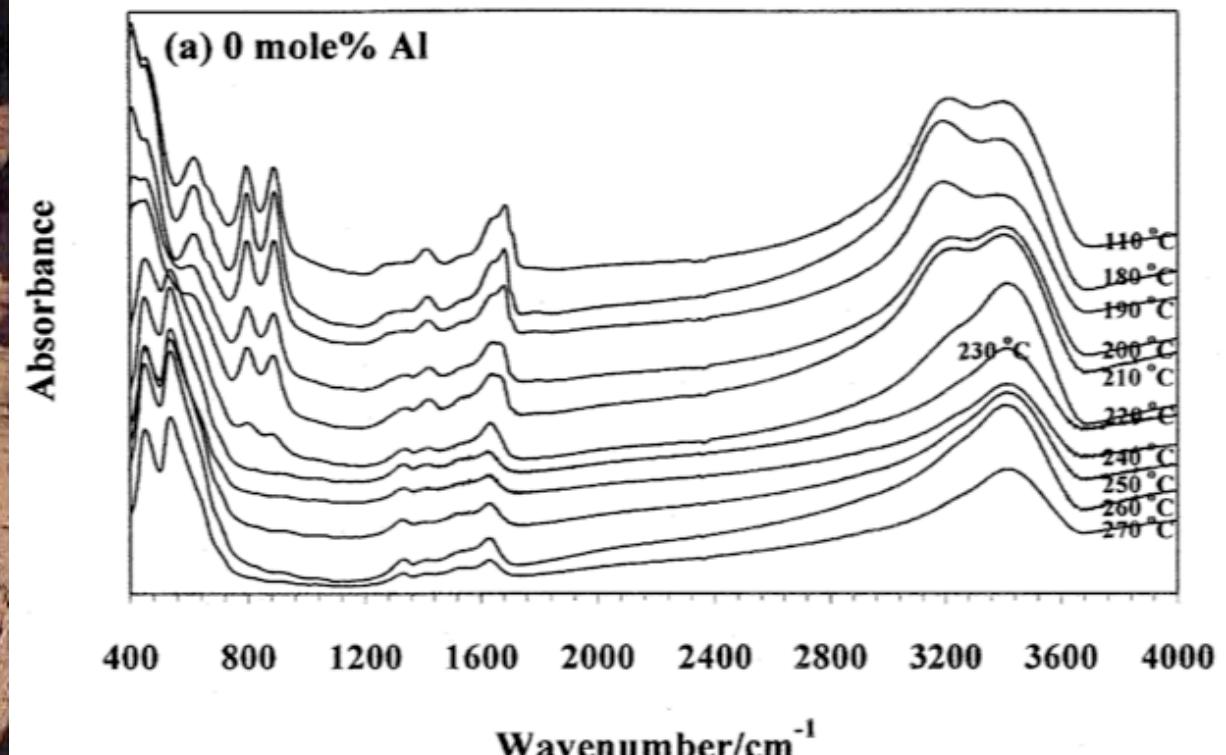
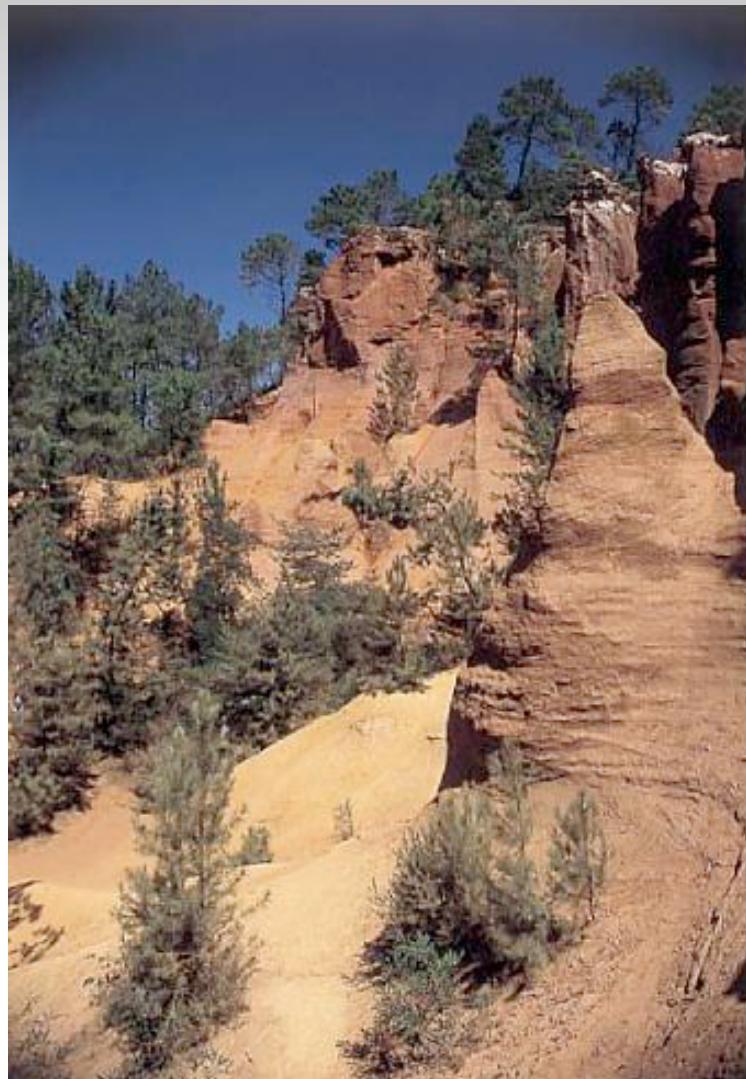
mineral transformations, but also for providing information about the

composition, the structural defects, the particle shape...



Theoretical modeling of IR spectra is useful

## Transformation of goethite to hematite



(Ruan et al. 2002, *Spectrochim Acta*)

Goethite-Hematite, Lubéron  
(Cornell and Schwertmann 2003)

## Al-bearing goethites

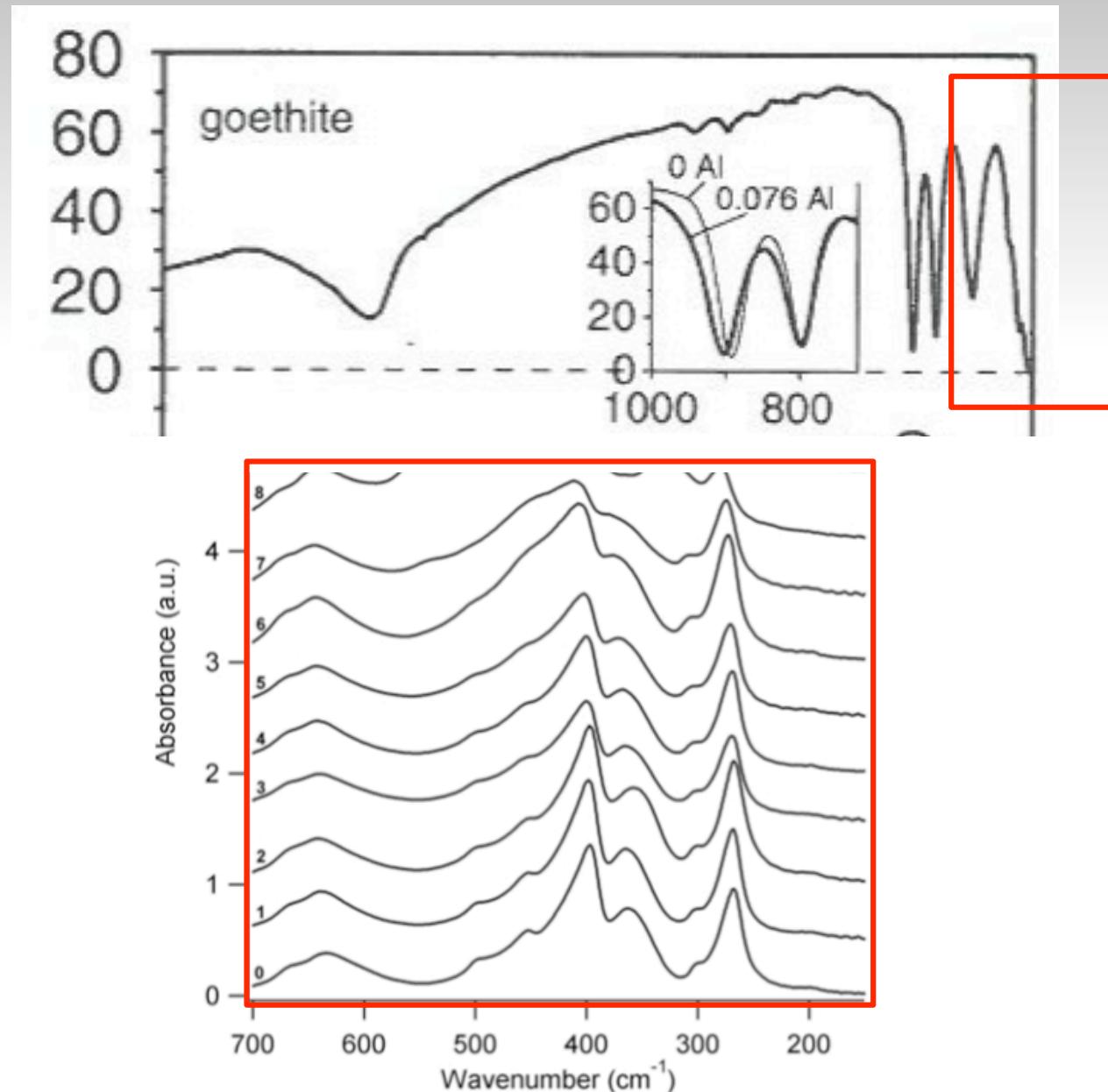
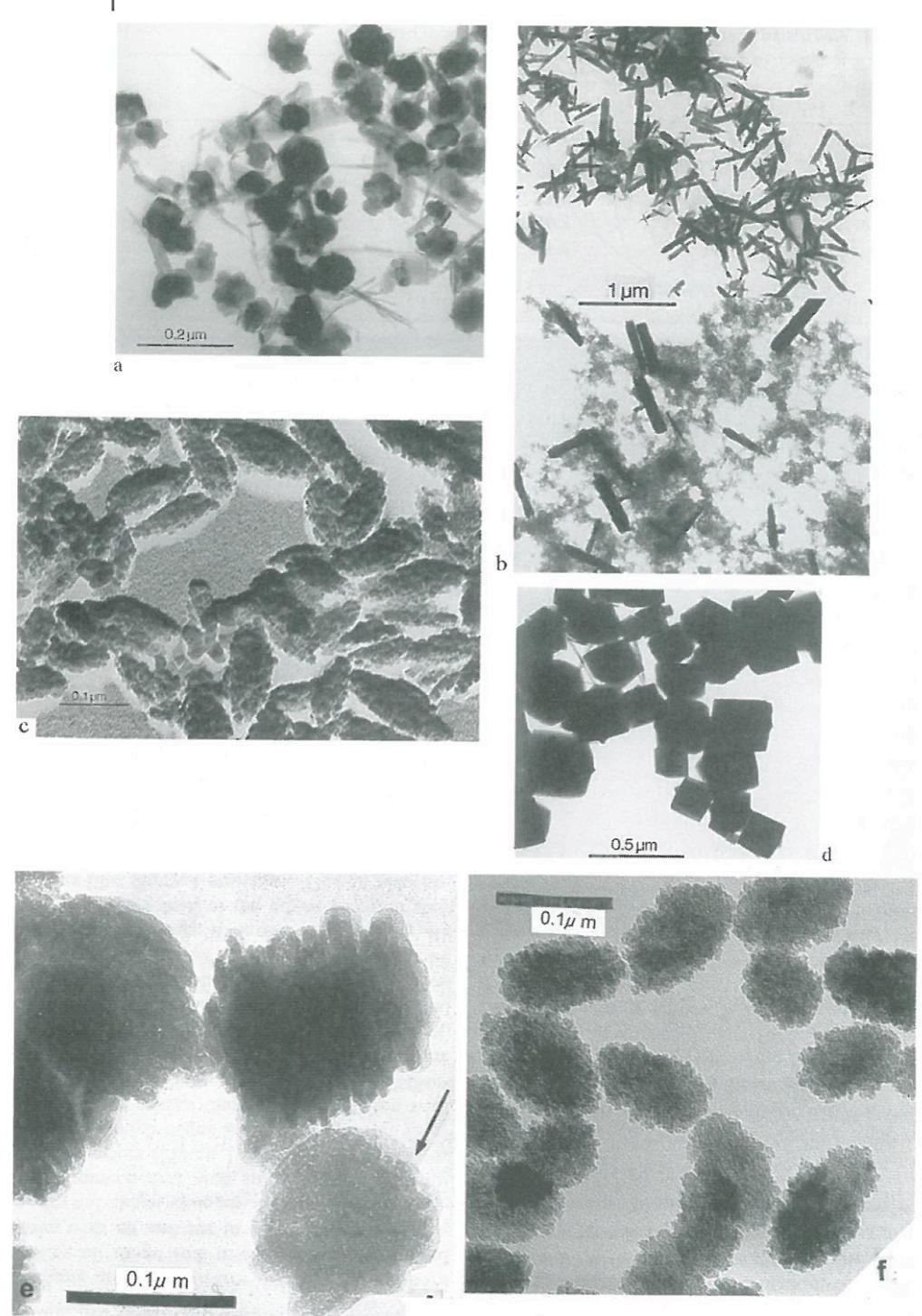


FIG. 4. IR spectra of Al-bearing goethites (see Table 2 for compositions of samples).  
(Blanch et al. 2008, *Miner Mag*)

# Hematite morphologies

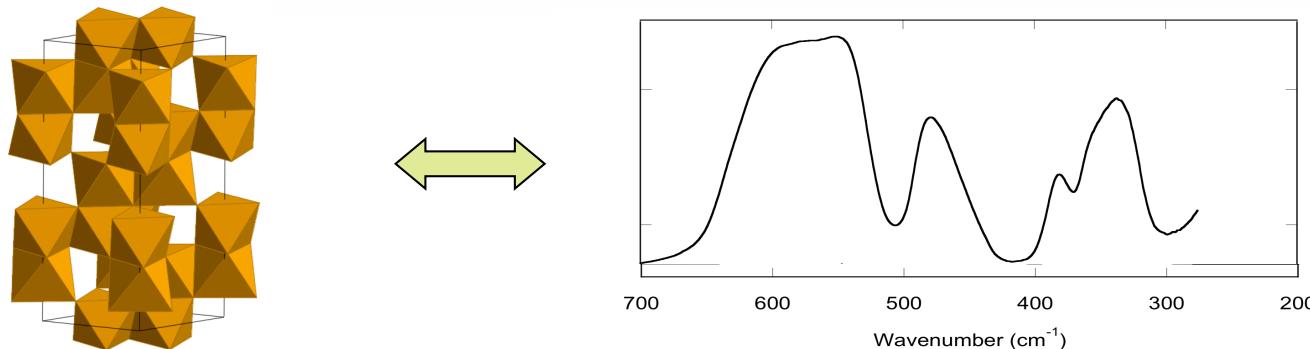
(Cornell and Schwertmann 2003)



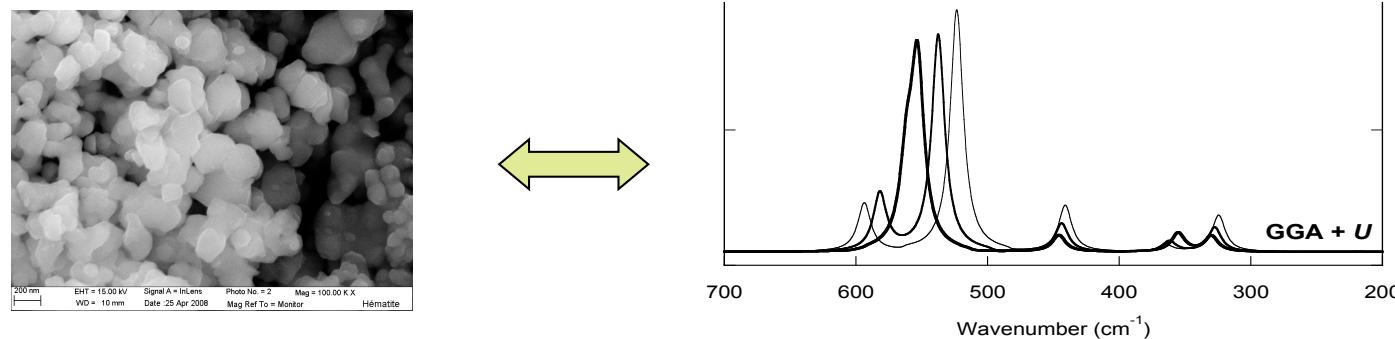
## Theoretical modeling of IR spectra

In the interpretation of IR spectra, first-principles calculations are useful for:

- finding unambiguous relations between IR spectrum, vibrational modes and crystal structure



- determining the contribution of actual features of mineral samples (e.g. particle shape, defects) to their IR spectrum



For a solid...

The induced polarisation is proportional to the macroscopic electric field existing in the material:

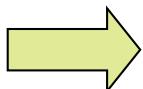
$$4\pi \mathbf{P}(\omega) = (\epsilon(\omega) - 1) \mathbf{E}(\omega) \quad \text{where } \epsilon(\omega) \text{ is the dielectric tensor}$$

*An approach in two steps :*

First-principles calculation  
DFT

Classical electromagnetic  
model

CRYSTAL



dielectric tensor  
 $\epsilon(\omega)$



IR SPECTRUM

- GGA (PBE)
- planewaves, pseudo-potentials
- PWscf, CASTEP

- experimental geometry
- sample micro-structure
- ...

# *Calculation of the low-frequency dielectric tensor $\epsilon(\omega)$*

$$S_{m,ij} = \left( \sum_{\tau,i'} Z^*_{ii',\tau} U_m(\tau,i') \right) \left( \sum_{\tau,j'} Z^*_{jj',\tau'} U_m(\tau',j') \right)$$

$\Gamma$  damping coefficient

$\epsilon(\infty)$  electronic dielectric tensor

## **Z\*** Born effective charge tensors

$U_m$  transverse normal modes

$\omega_m$  transverse vibrational frequencies

} calculated using DFPT: second derivatives  
of total energy w.r.t. atomic displacements  
and/or macroscopic electric field

- harmonic approximation, athermal limit
  - vibrational modes at the center of Brillouin zone
  - accuracy: effective charges  $\sim 10\%$ ; frequencies  $\sim 3\%$ ;
  - contains the microscopic information

(Gonze and Lee 1997; Baroni et al. 2001)

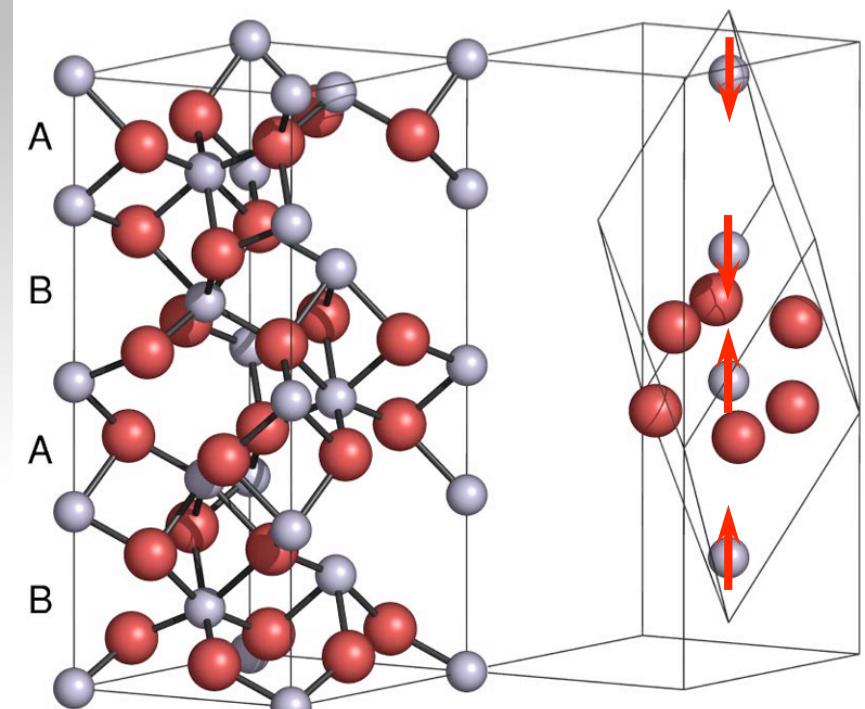
## Hematite ( $\alpha\text{-Fe}_2\text{O}_3$ )

Rhombohedral symmetry

Antiferromagnetic ( $T < 260$  K)



Highway on a hematite-coloured soil,  
Bolivia (Cornell and Schwertmann 2003)



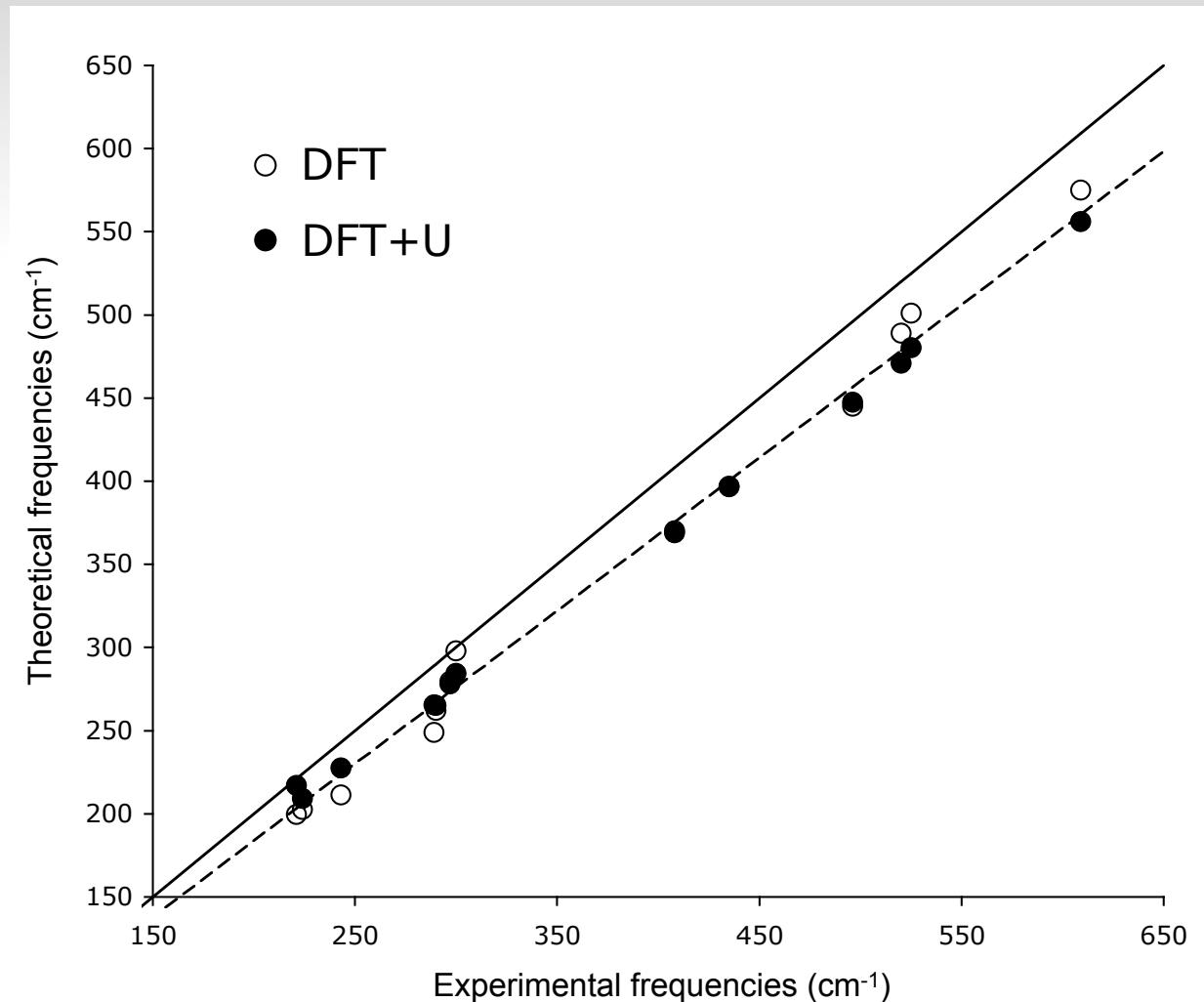
DFT+ $U$  needed to describe correctly electronic  
and magnetic properties

GGA (PBE) +  $U$  (3.3 eV)  
Ultrasoft pseudopotentials  
Wavefunctions cutoff = 40 Ry  
4x4x4  $k$ -points



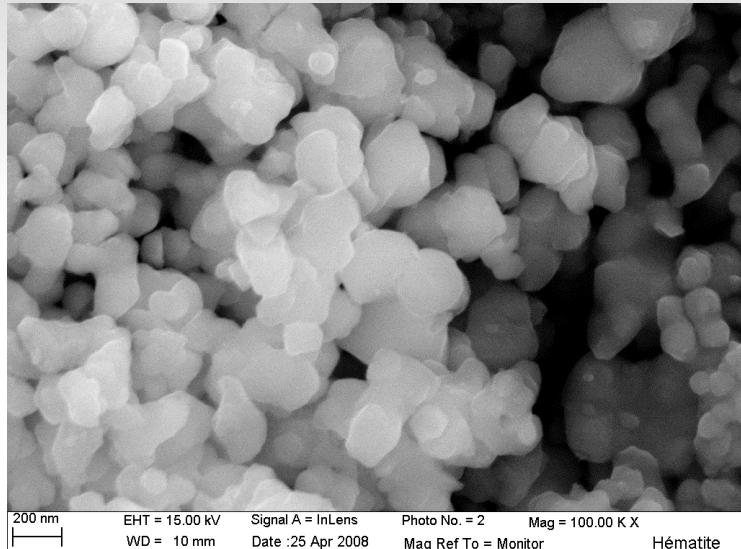
	GGA	GGA+ $U$	Exp
$a_{\text{hex}}$	5.023	5.098	5.035
$c_{\text{hex}}$	13.901	13.946	13.747
$c_{\text{hex}}/a_{\text{hex}}$	2.77	2.74	2.73
mag. mom. ( $\mu_{\text{B}}$ )	3.87	4.19	4.6-4.9
Band gap (eV)	0.7	1.5	2.0

## Theoretical vs. experimental frequencies of transverse optical modes

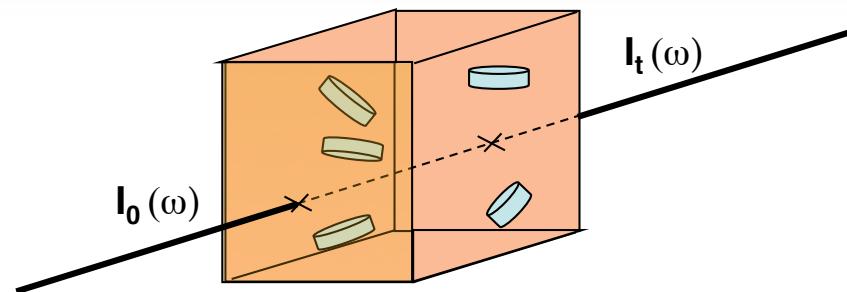


Blanchard et al. (2008, Am Miner)

## *Modeling the powder absorption spectra*



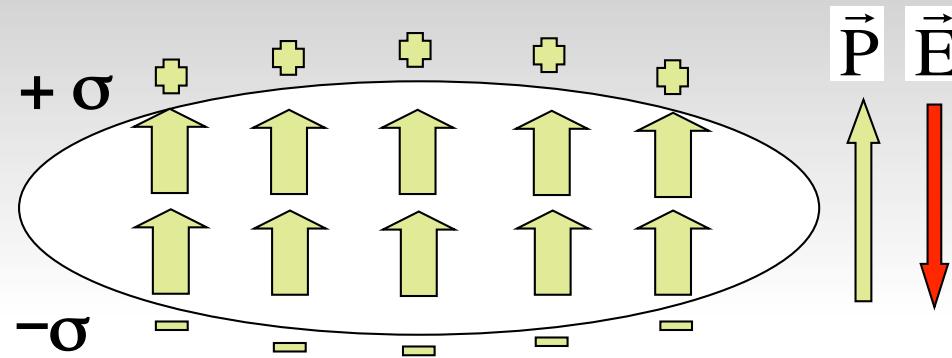
Less than 1 wt. % of powder sample  
is diluted in a non-absorbing matrix



Model:

- ellipsoidal particles
- isolated particles in homogeneous external medium (KBr)
- particle size  $< \lambda_{\text{IR}}$  (0.74 – 300  $\mu\text{m}$ )

IR absorption depends on the internal macroscopic electric field  
( $\neq$  applied electric field)



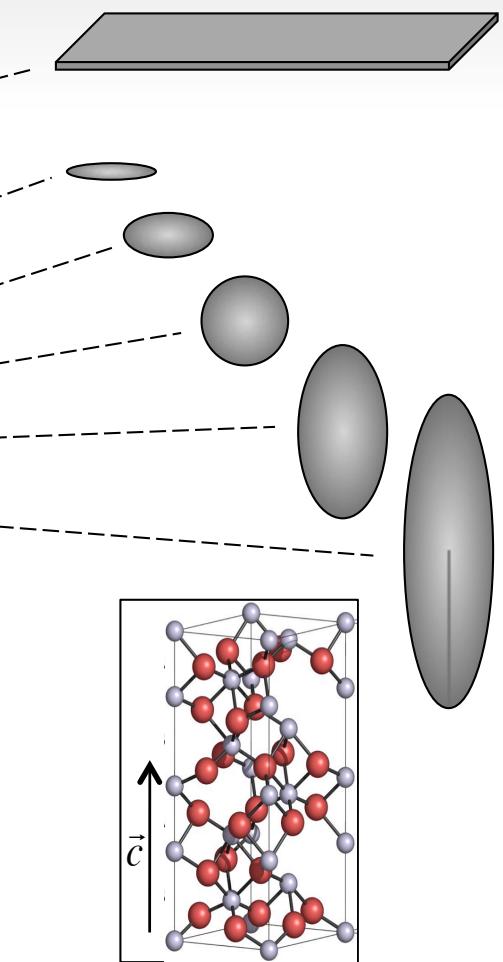
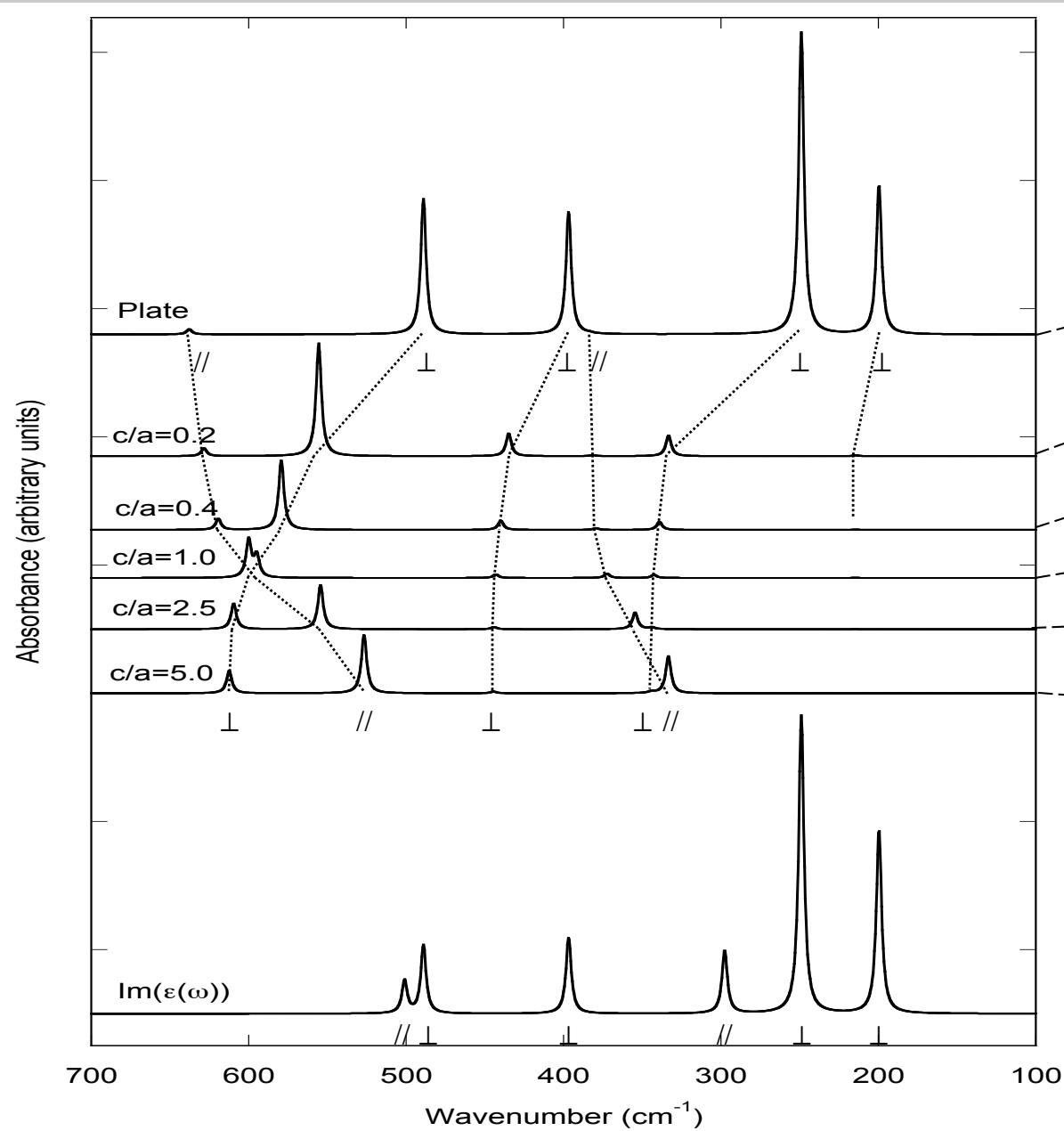
The electric field induced by surface charges in the polarized particles shifts the absorption bands and affects their intensity

$A(\omega)$  proportional to the orientational and temporal average of the electromagnetic power dissipated in the particles ( $W(\omega)$ ).

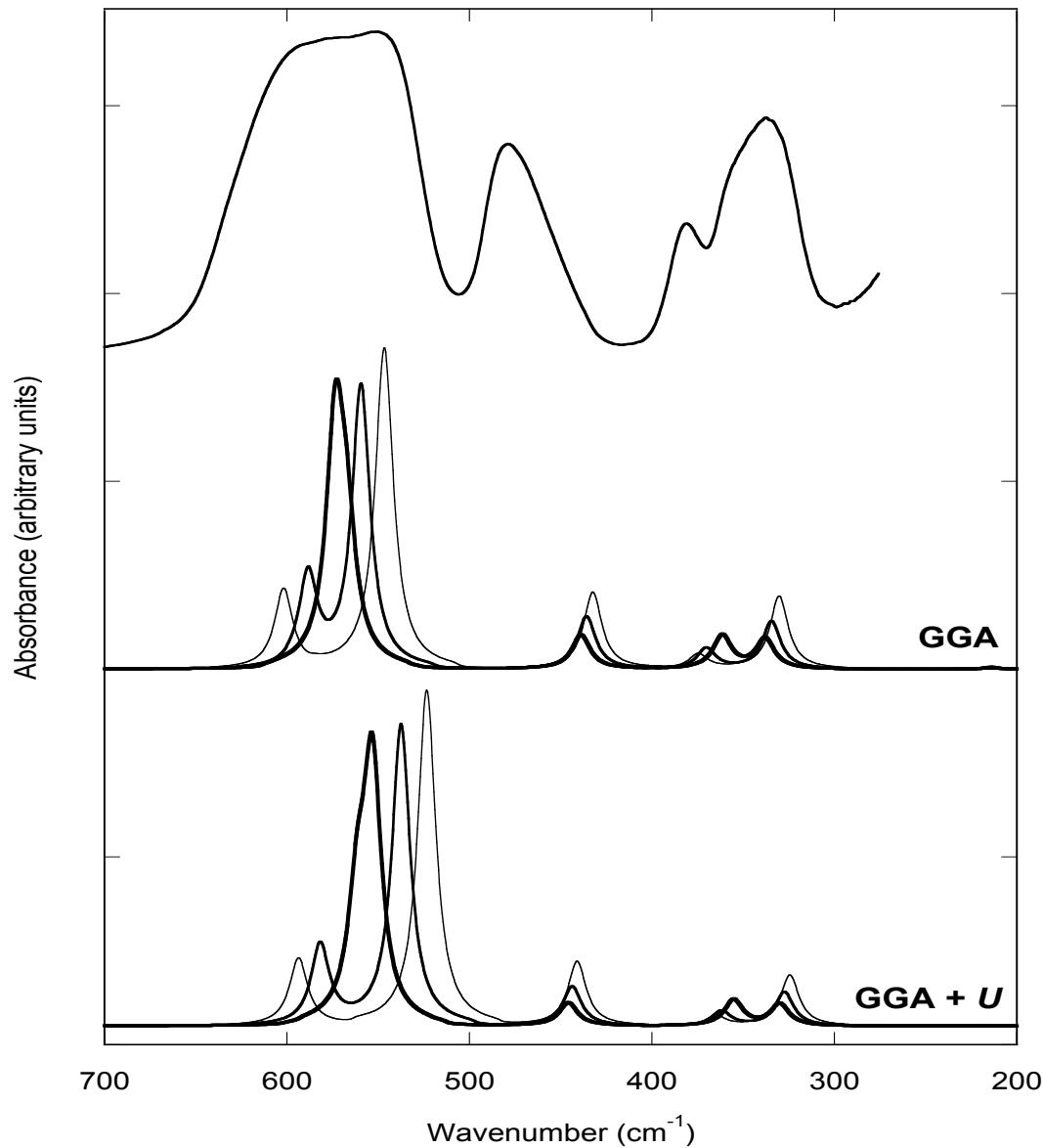
$W(\omega)$  : function of the external electric field  $E_{\text{KBr}}$ , and dielectric tensors of the sample  $\epsilon(\omega)$  and external medium  $\epsilon_{\text{KBr}}$

$$\rightarrow W(\omega) \propto |E_{\text{KBr}}|^2 \omega \operatorname{Im}[\epsilon(\omega) - \frac{\mathbf{n} \cdot (\epsilon(\omega) - \epsilon_{\text{KBr}})^2 \cdot \mathbf{n}}{(3 \mathbf{n} \cdot \epsilon(\omega) \cdot \mathbf{n})}]$$

## *Effects related to particle shape in the IR spectrum of hematite*



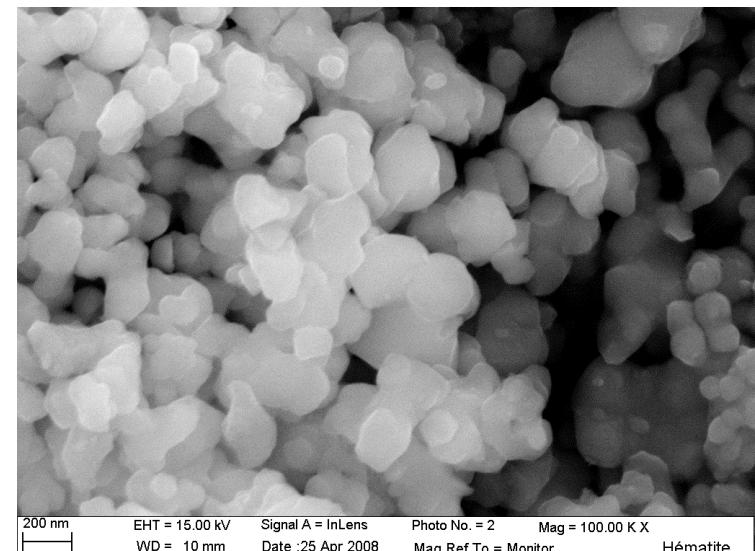
Blanchard et al. (2008, Am Miner)



Broadening related to non-ellipsoidal particles and distribution of particle shapes.

Experimental (top) and theoretical (bottom) IR powder absorption spectra of hematite

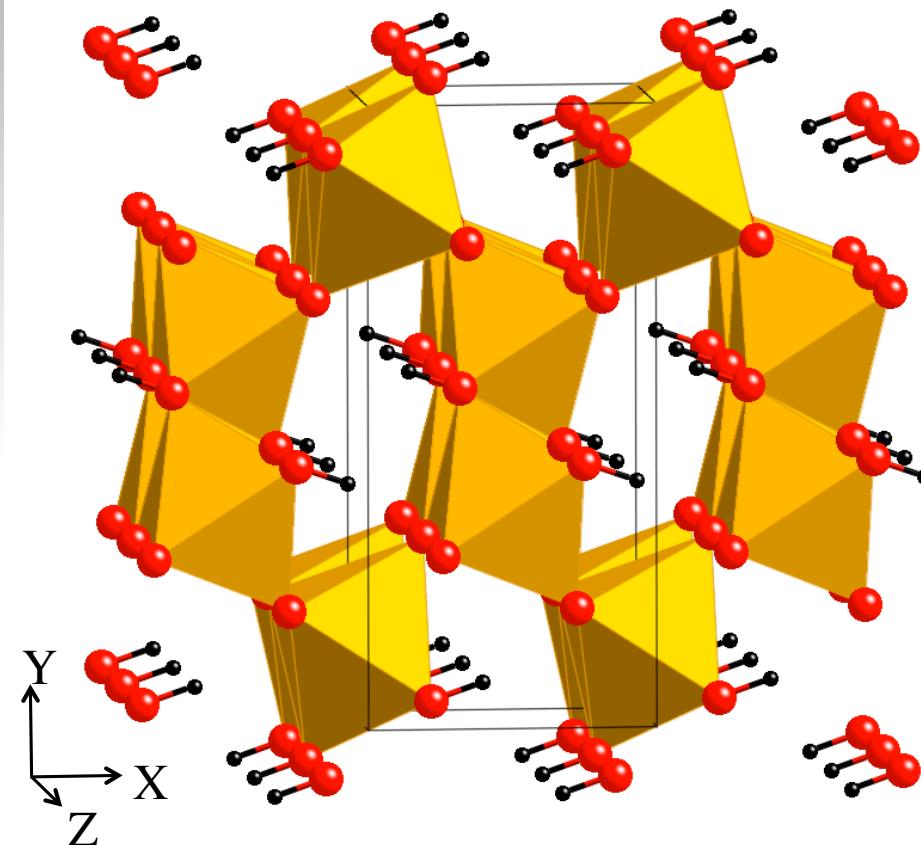
$c/a$        $c/a$   
0.4      1.0  
○      ●



## Goethite ( $\alpha\text{-FeOOH}$ )

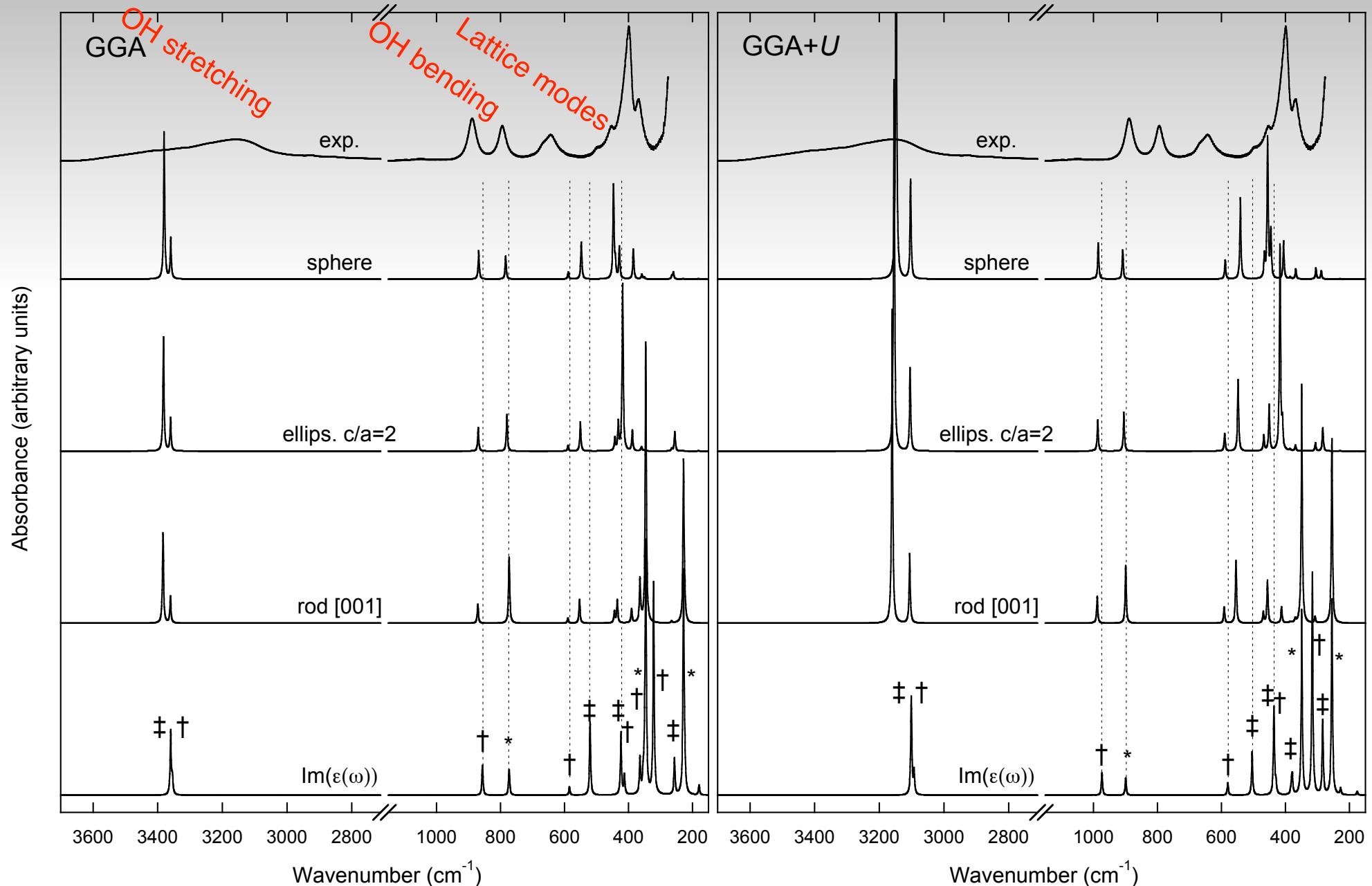
Orthorhombic unitcell ( $Pbnm$ )

GGA (PBE) +  $U$  (3.3 eV)  
Ultrasoft pseudopotentials  
Wavefunctions cutoff = 40 Ry  
 $4 \times 2 \times 6$   $k$ -points



	a	b	c
Goeth. GGA	4.652	10.013	3.026
Goeth. GGA+ $U$	4.638	10.085	3.067
Exp. (Yang et al. 2006)	4.598	9.951	3.018

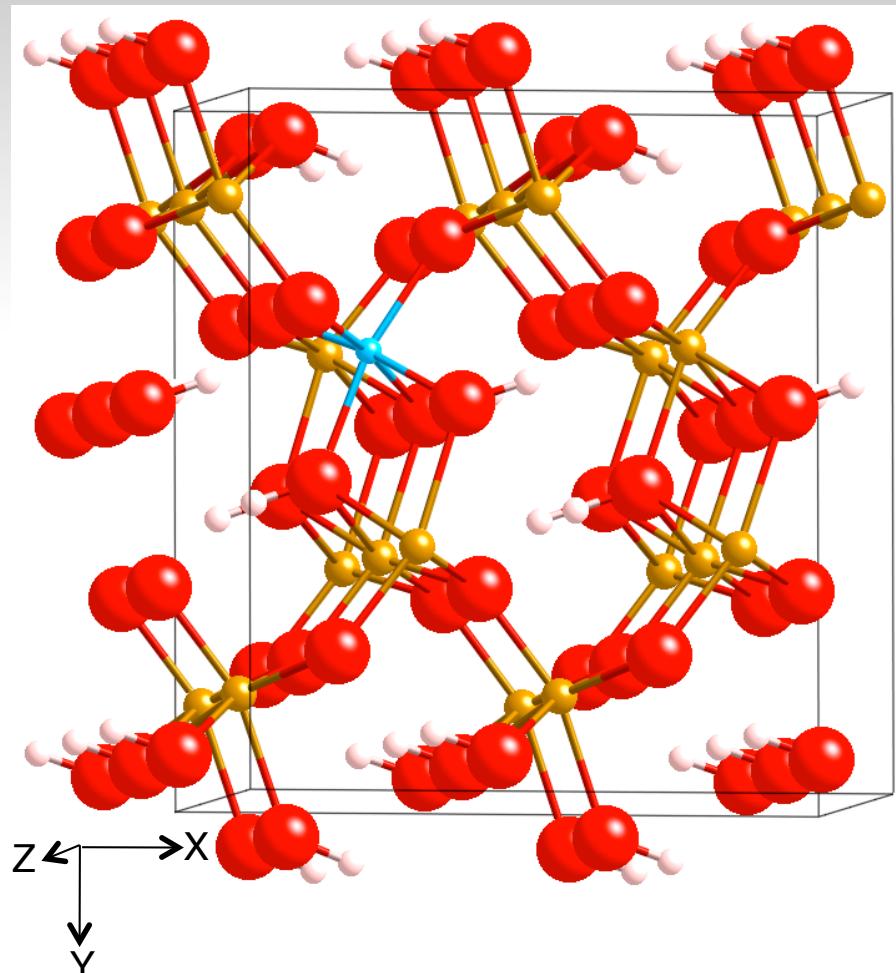
# Experimental (top) and theoretical (bottom) IR powder absorption spectra of goethite



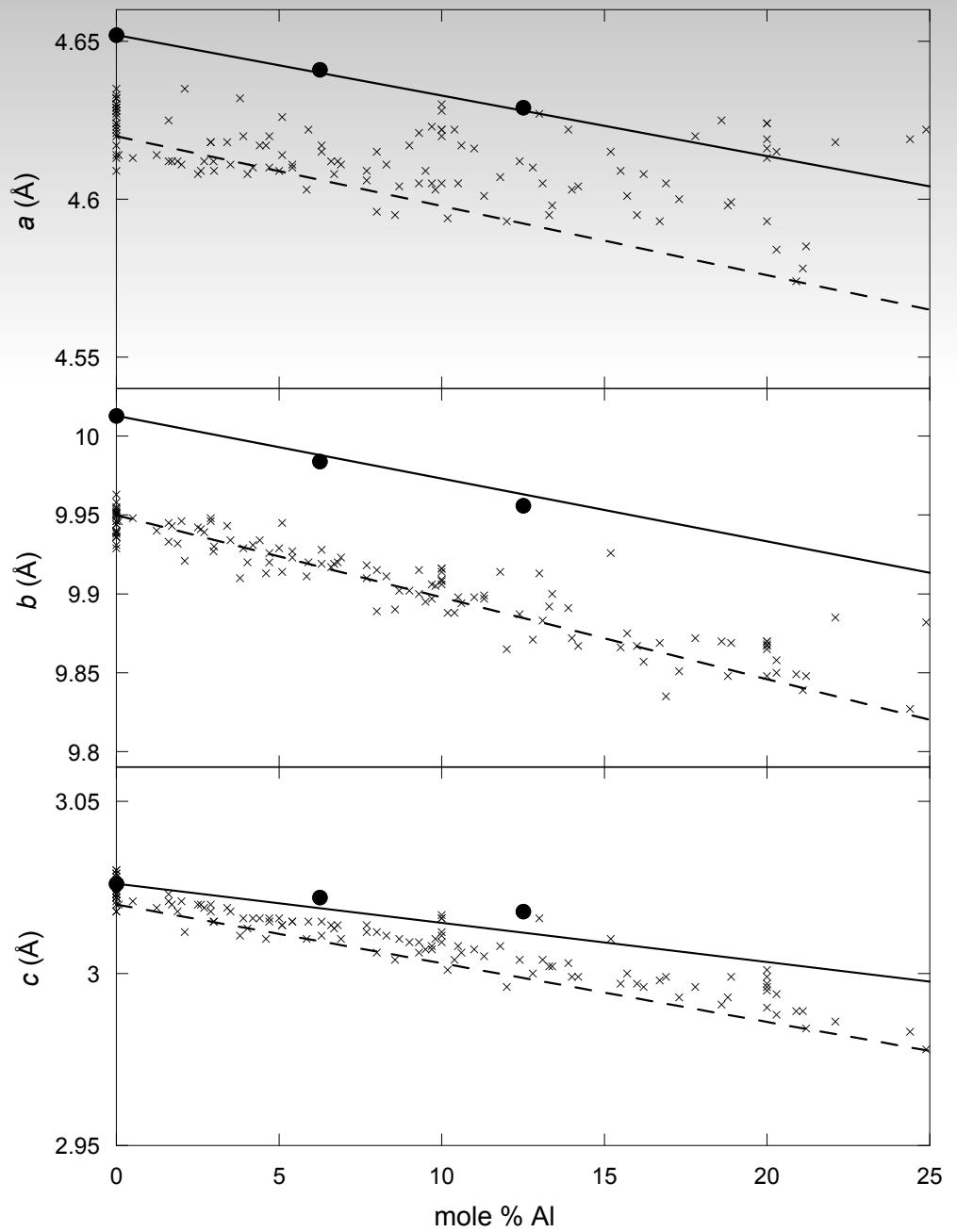
$\ast$ ,  $\dagger$  and  $\ddagger$  for  $B_{1u}$ ,  $B_{2u}$  and  $B_{3u}$  symmetries, respectively

Blanchard et al. (PCM, under review)

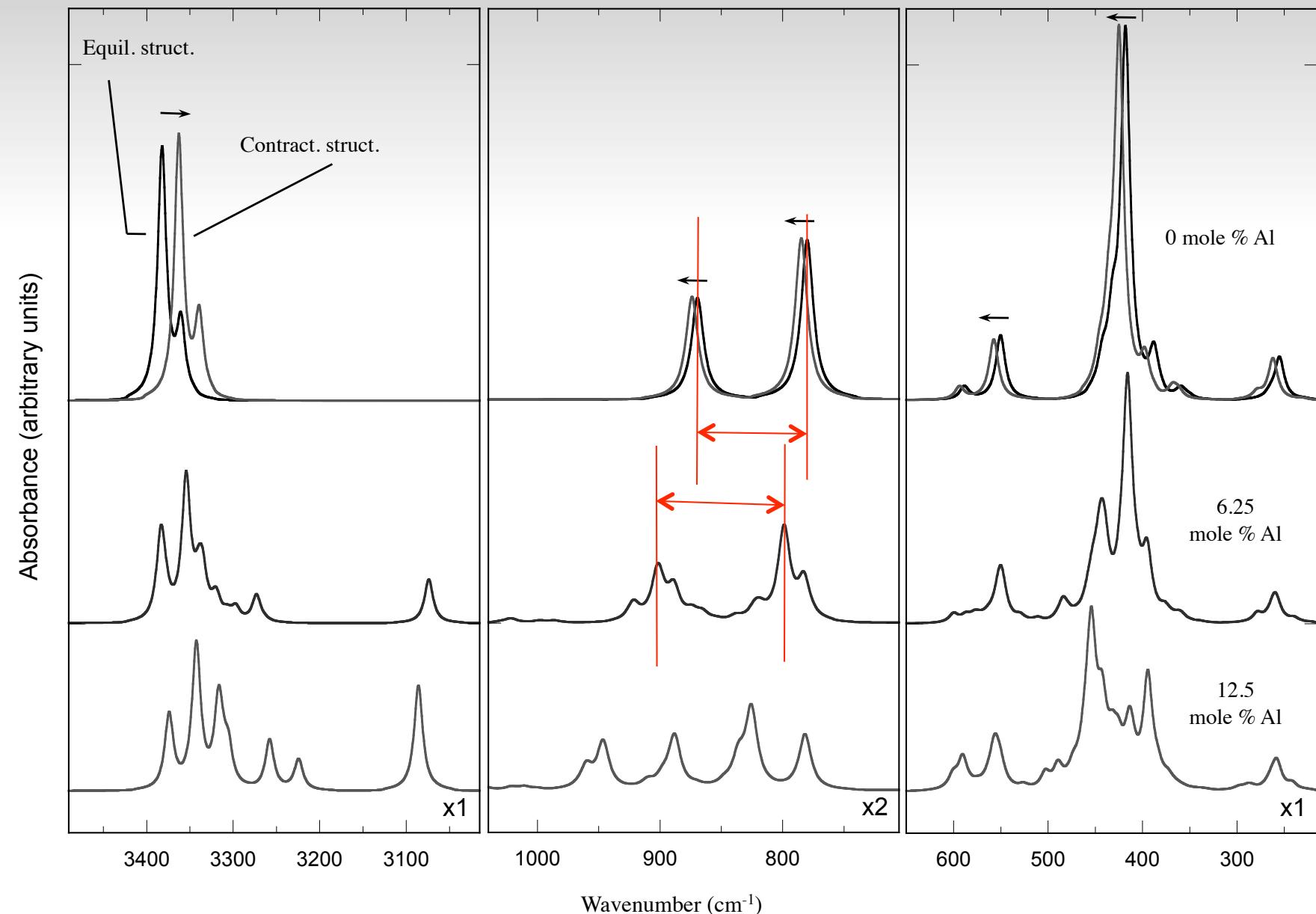
# Experimental and theoretical (GGA) lattice parameters of goethite versus Al content



The  $\text{AlO}_6$  octahedron in goethite is sensibly the same as in diasporite  
(i.e. almost full relaxation)



# Effect of the Al substitution on the theoretical (GGA) IR absorption spectrum of goethite

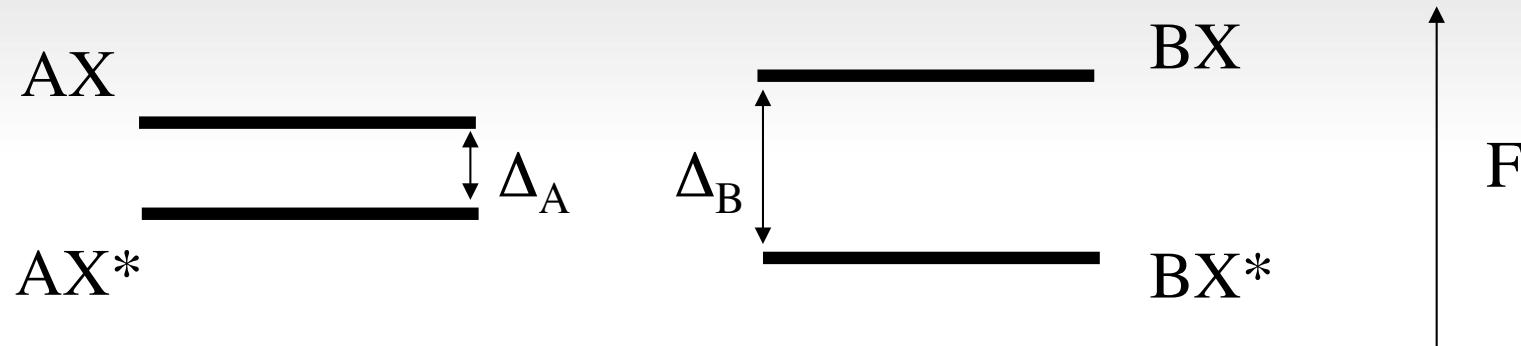


## *Conclusions I*

- First-principles methods accurate enough to enable a direct comparison with experiment
- A single theoretical framework to interpret the results of various spectroscopic measurements (IR, Raman, NMR, XAS, ...) on complex samples
- Microscopic vs. macroscopic effects in IR powder spectra

## Theoretical isotopic fractionation

X isotope exchange between two phases A and B:



$$\ln (K) = (\Delta_A - \Delta_B) / kT$$

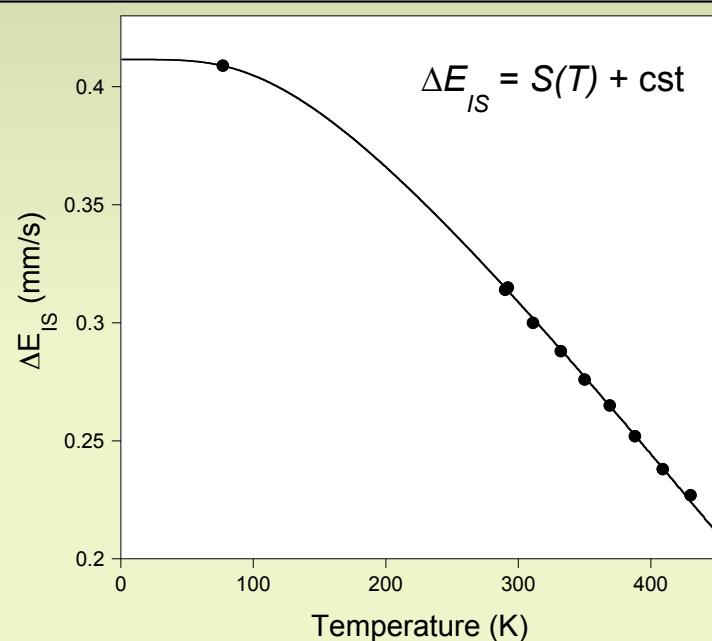
$$F = -kT \ln (Q) \Rightarrow K = (Q_{AX^*}/Q_{AX})/(Q_{BX^*}/Q_{BX})$$

For a crystal, the partition function is: 
$$Q = \left[ \prod_{i=1}^{3Nat} \prod_{\{q\}} \frac{e^{-hv_{q,i}/(2kT)}}{1 - e^{-hv_{q,i}/(kT)}} \right]^{1/N_q}$$

$\beta$ -factor = equilibrium isotopic fractionation factor between phase A and an ideal gas of X

From 1<sup>st</sup> order thermodynamic perturbation theory:

$$\ln \beta = \frac{m^* - m}{m} \left( \frac{K}{RT} - \frac{3}{2} \right)$$



In Mössbauer, the kinetic energy of the iron sublattice is related to the second-order Doppler shift ( $S(T)$ )

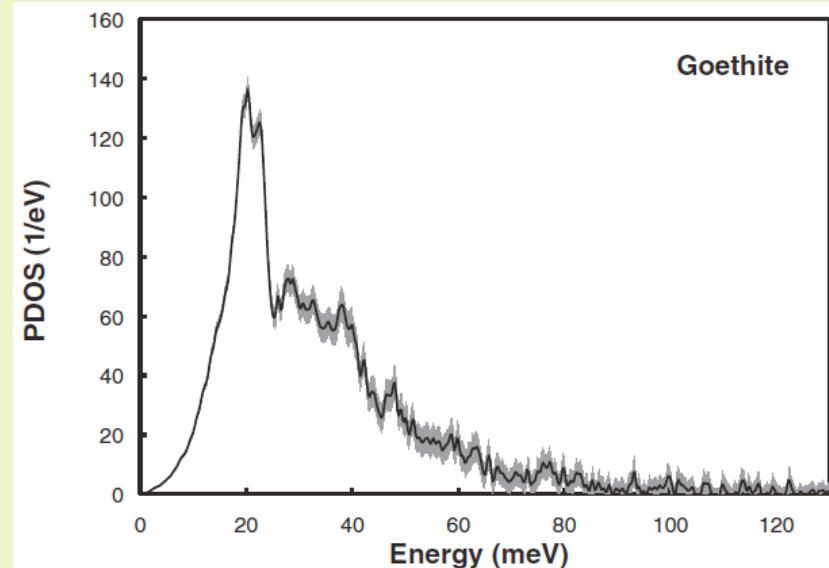
$$S(T) = -\frac{K(T)_{^{57}Fe}}{m c}$$

velocity of light

In Nuclear Resonant Inelastic X-ray Scattering (NRIIXS), the kinetic energy is related to the partial vibrational density of state (PDOS)

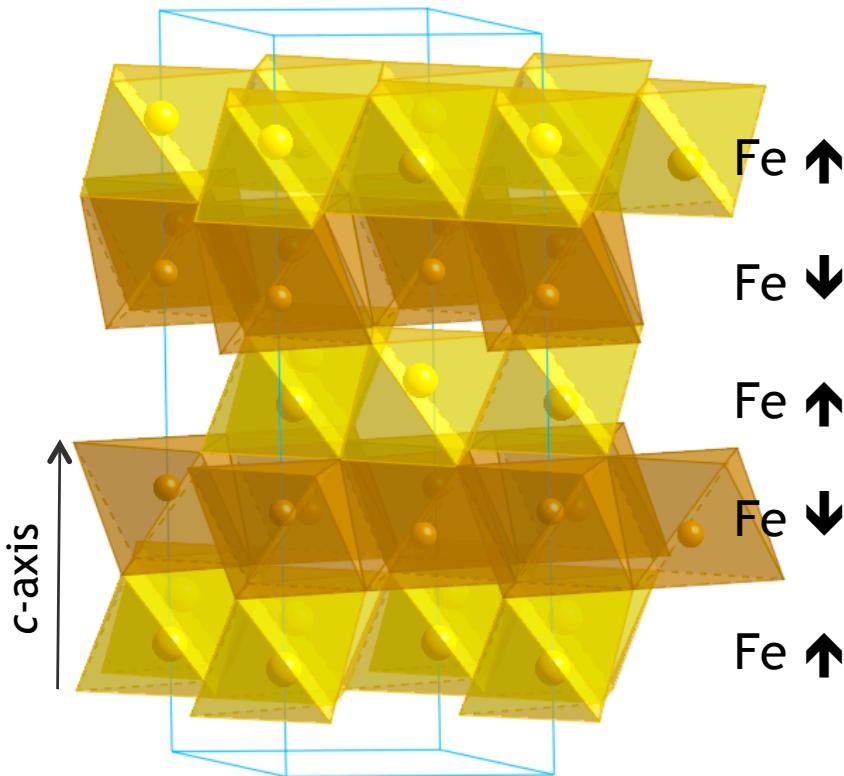
$$K_{^{57}Fe} = \frac{3}{2} \int_0^{e_{\max}} E(e, T) g(e) de$$

↑ Einstein function      ↑ PDOS norm. to 1



## Computational method

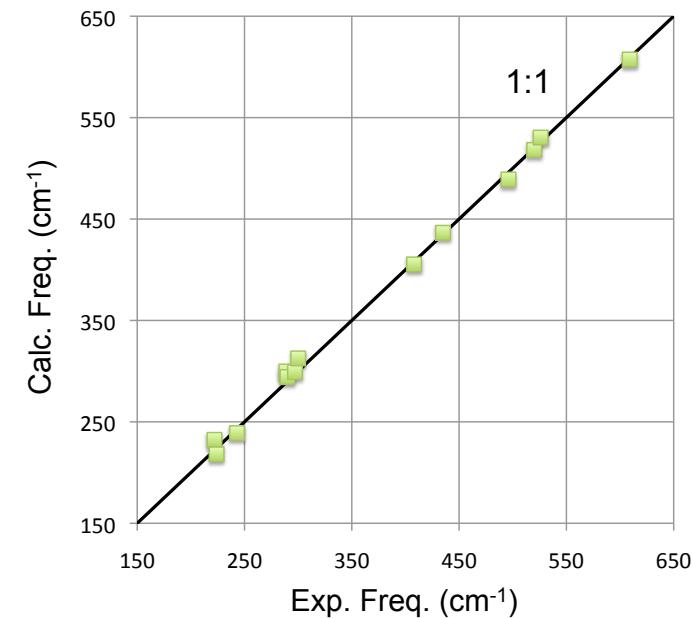
- Density Functional Theory  
(Plane-waves / Ultrasoft pseudopotentials /  
Experimental cell parameters / GGA+U )
- Hematite ( $\alpha\text{-Fe}_2\text{O}_3$ ), hexagonal sym.,  
antiferromagnetic



Atomic positions

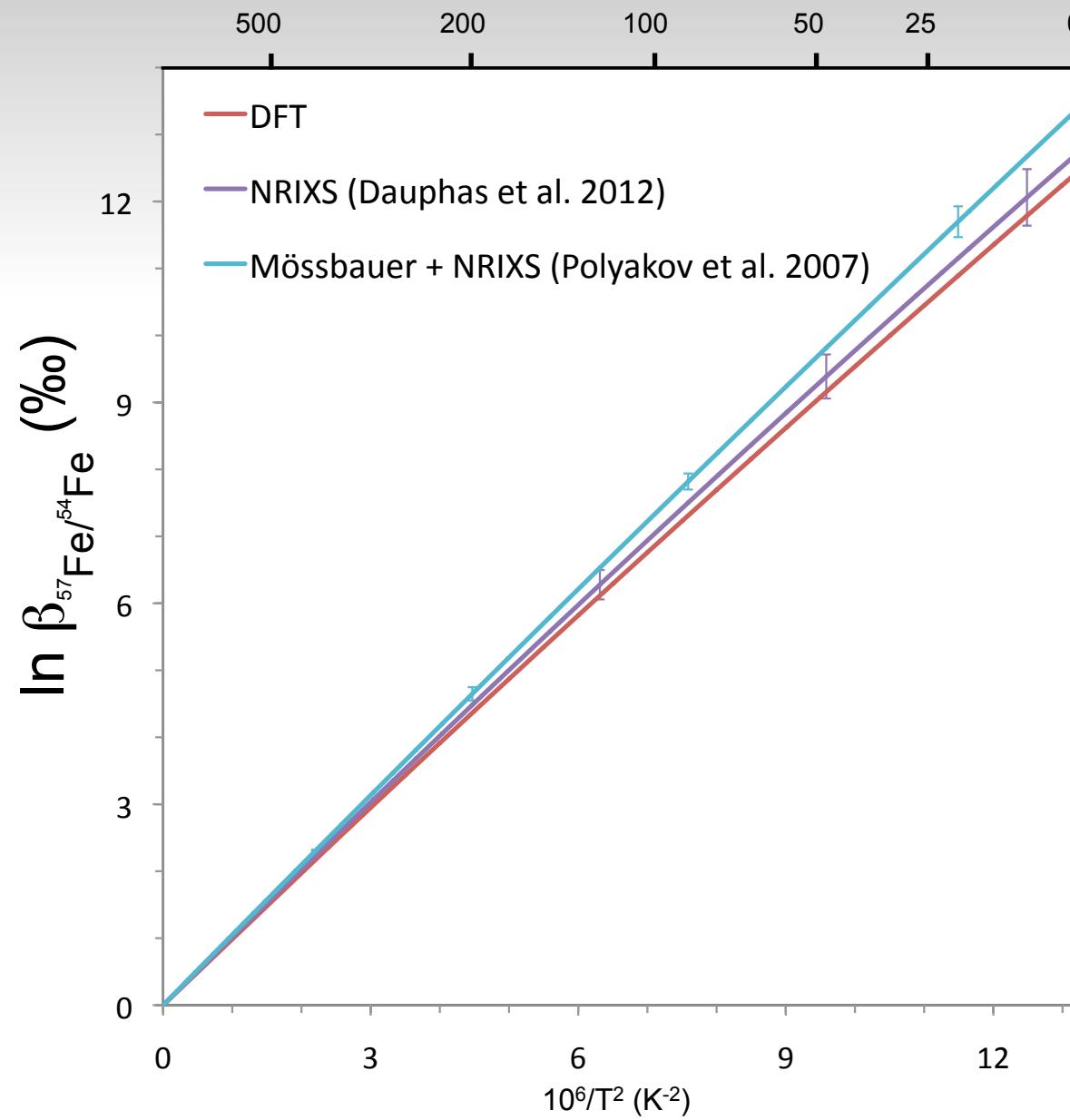
	$x_{\text{Fe}}$	$x_{\text{O}}$
Calc.	0.1449	-0.0558
Exp.	0.1447	-0.0556

Raman and IR frequencies

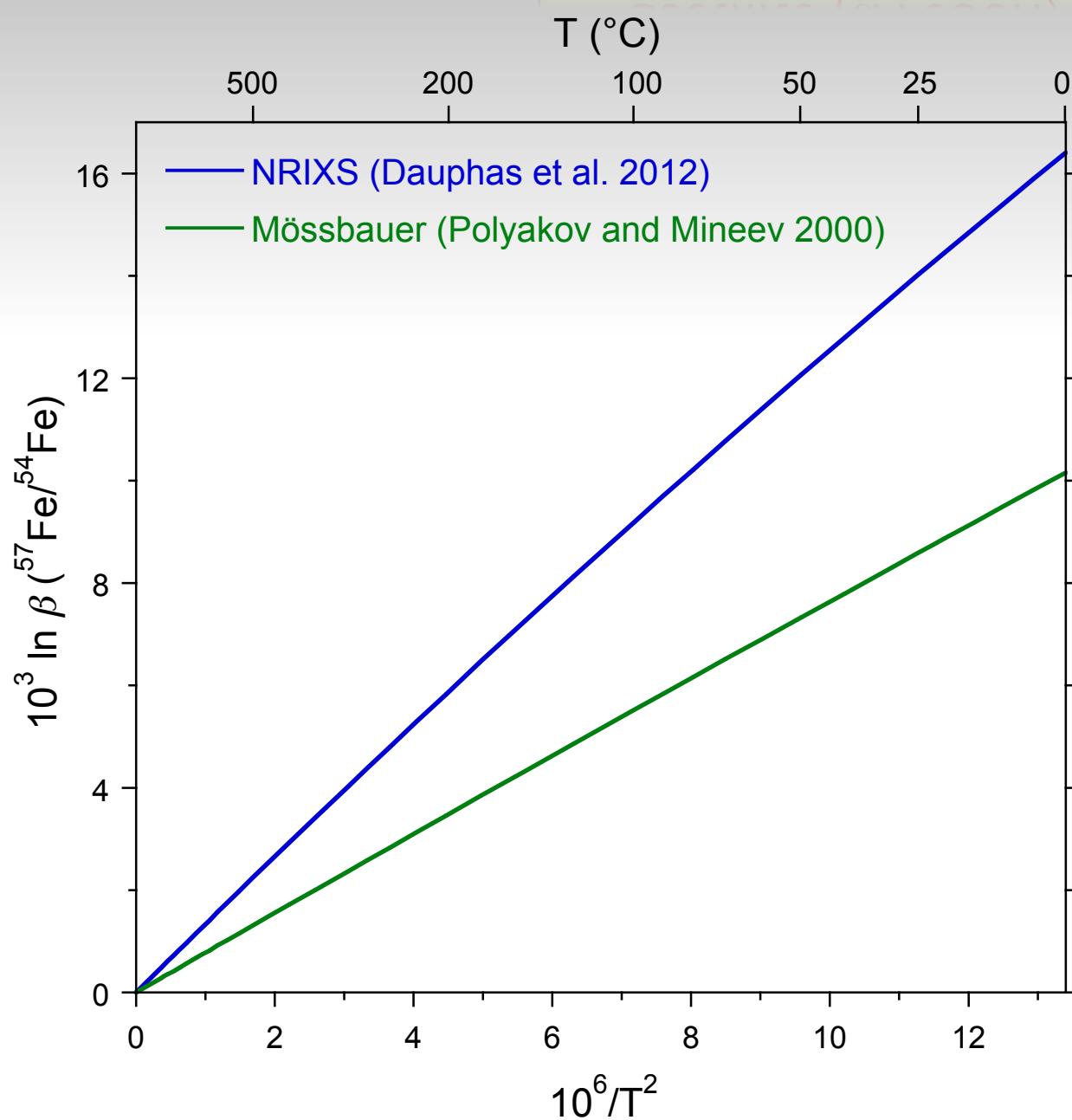


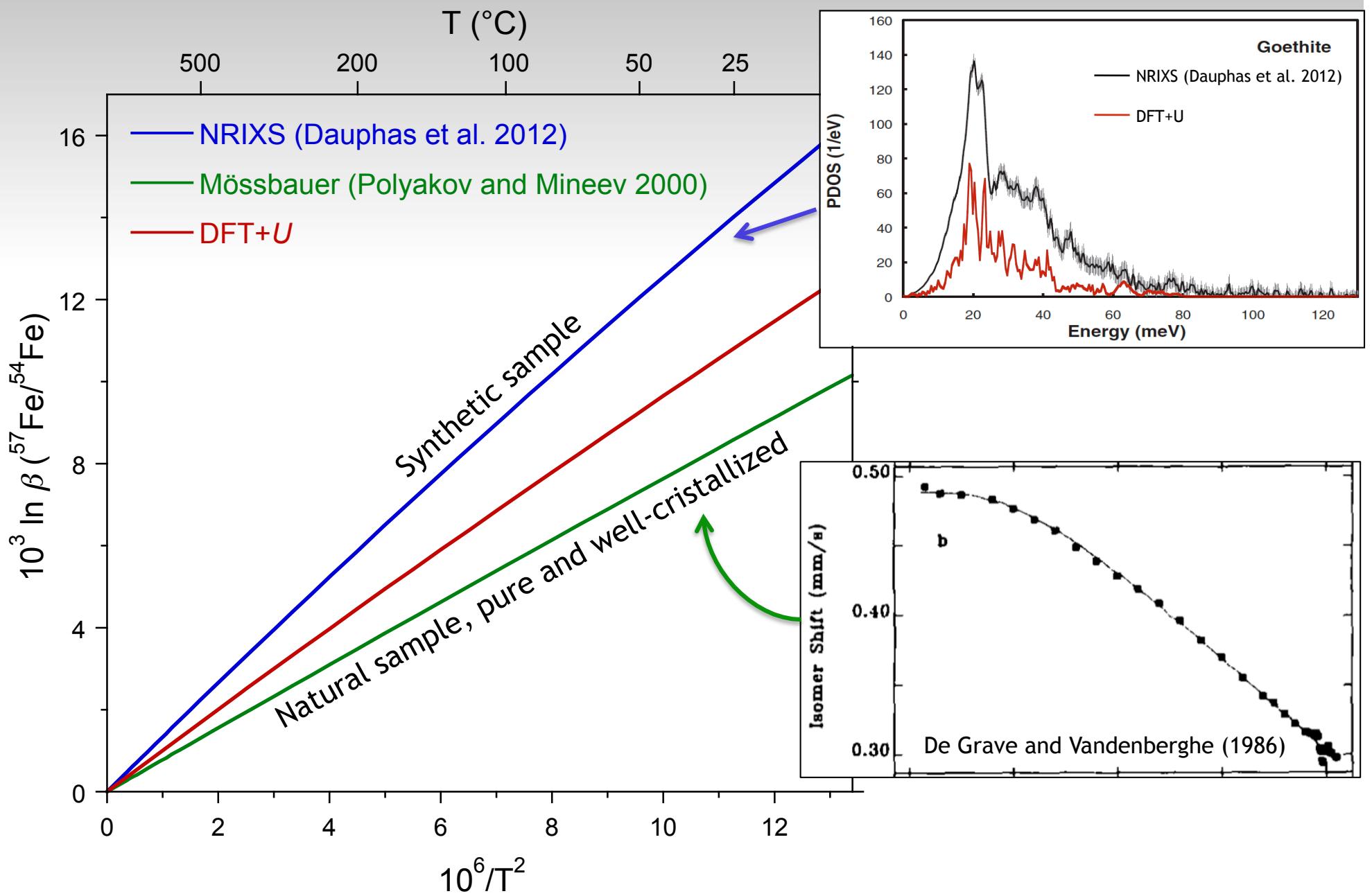
## Hematite ( $\alpha\text{-Fe}_2\text{O}_3$ )

T (°C)



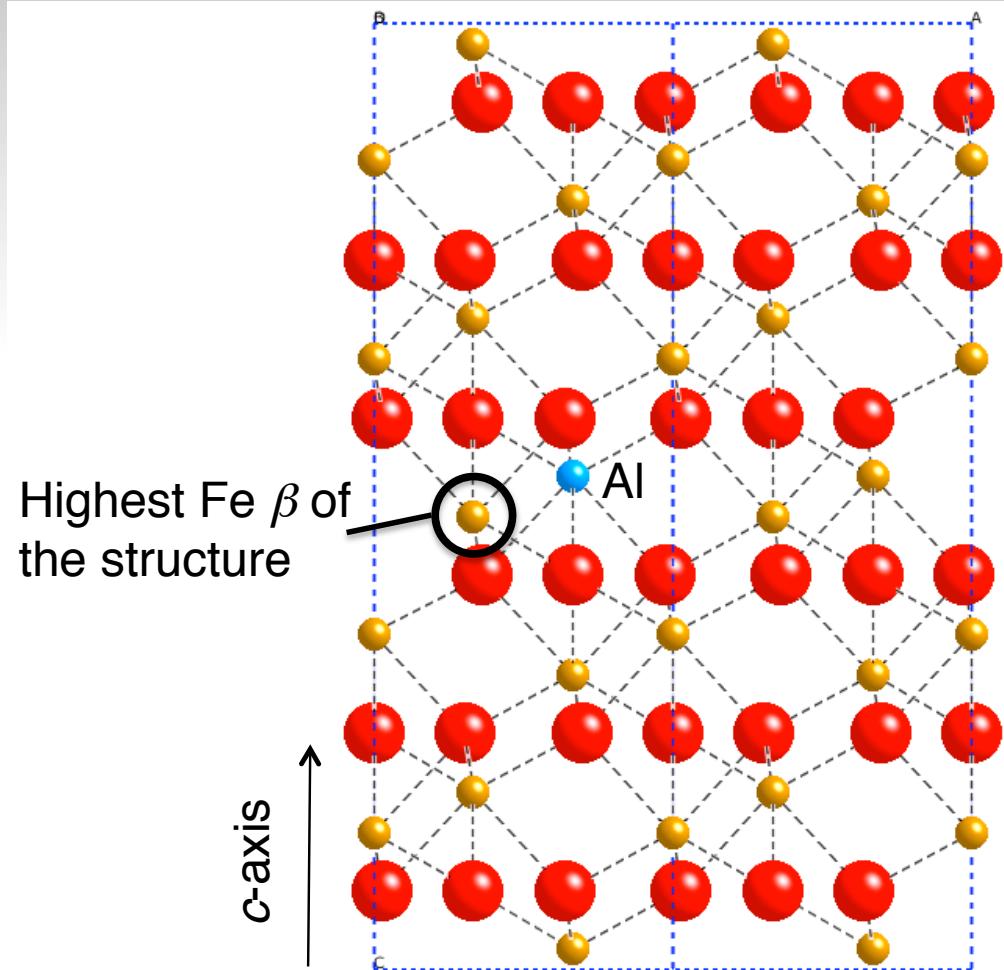
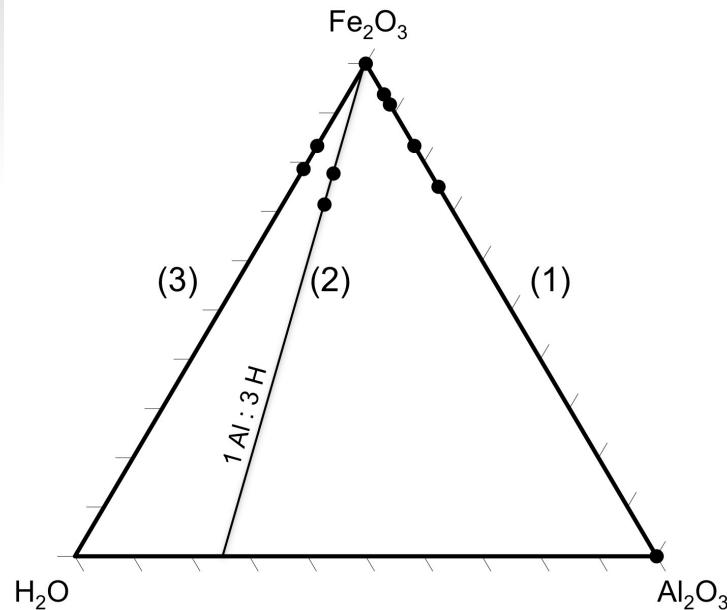
## Goethite ( $\alpha\text{-FeOOH}$ )





## *Effect of cationic substitution (Al, H)*

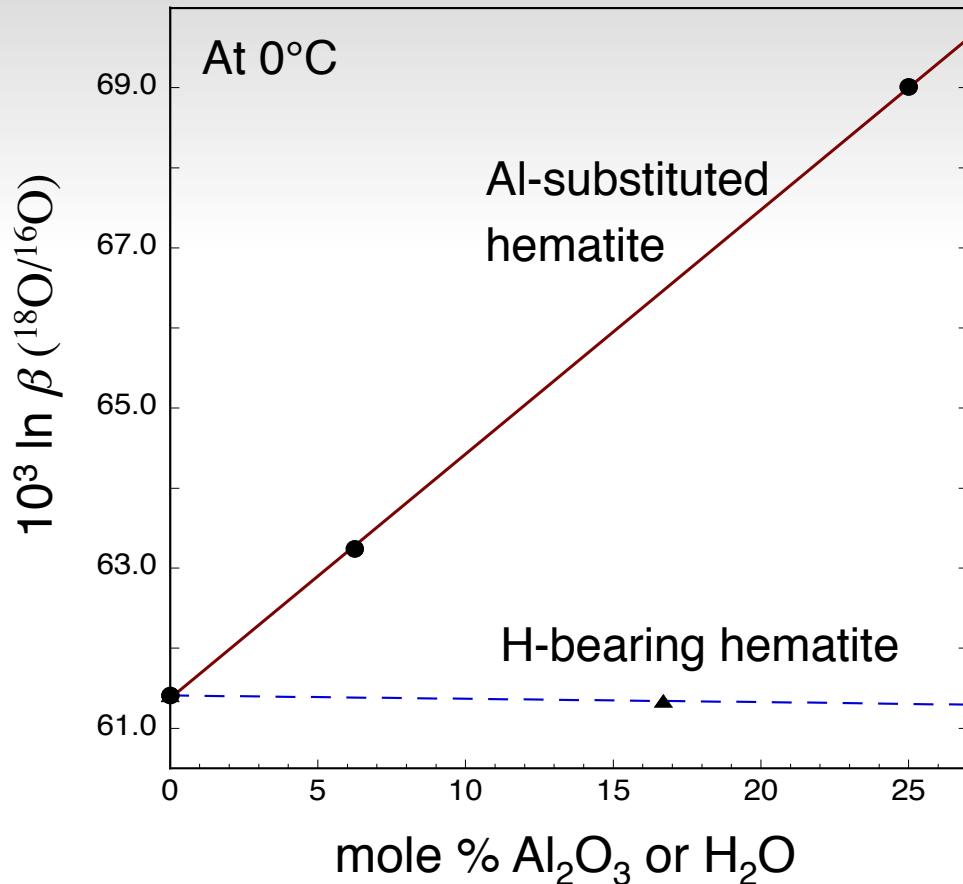
Compositions investigated



The iron  $\beta$  is higher for Fe atoms located next to the Al impurity.

The effect is larger when Al is in the edge-sharing octahedron rather than in the corner-sharing or face-sharing octahedron.

## Oxygen $\beta$ -factor of hematite vs Al or H content



- An incorporation of 18 mole %  $\text{Al}_2\text{O}_3$  in hematite would increase the oxygen  $\beta$  of  $\sim 5.5\text{\%}$  at  $0^\circ\text{C}$ . This effect is sufficiently large to be measurable and to affect the interpretation of natural isotopic compositions.
- The effect of H incorporation is found to be negligible. This is explained by a complex local relaxation of the defective structure leading to a compensation of local  $\beta$  values.

## Conclusions II

- Comparing the data obtained from the different techniques (DFT, Mössbauer, NRIXS) enables to obtain reliable equilibrium isotopic fractionation factors.
- First-principles calculations also allow to investigate the mechanisms controlling isotopic fractionations at the molecular scale (crystallographic sites, surface effect, adsorption processes ...)

## Acknowledgements

### Collaborators

M. Lazzeri, F. Mauri, G. Morin, P. Giura, K. Béneut, M. Guillaumet, H. Yi,  
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A. Floris (*King's College London*)



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